



RBRC workshop on  
“High  $p_T$  Physics at RHIC”

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# Spin Physics

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## ● A Brief History of Spin

## ● Spin Physics at RHIC

- ▷ theoretical framework of perturbative QCD
- ▷ accessing the gluon polarization
- ▷ transverse spin phenomena

## ● Summary & Outlook

# I. A Brief History of Spin

- the characteristic property of a particle besides mass
- determines whether the particle is a fermion or boson
- practical effects prevail the whole of science:
  - structure of composite objects
  - material properties like superconductivity
  - gateway to new physical phenomena like BEC



A world without spin would be boring!

The story of the proton spin begins in 1927:

Hund: rot. part of specific heat of  $H_2$  molecule

Hori: band spectrum of  $H_2$

↔ proton's spin & statistics

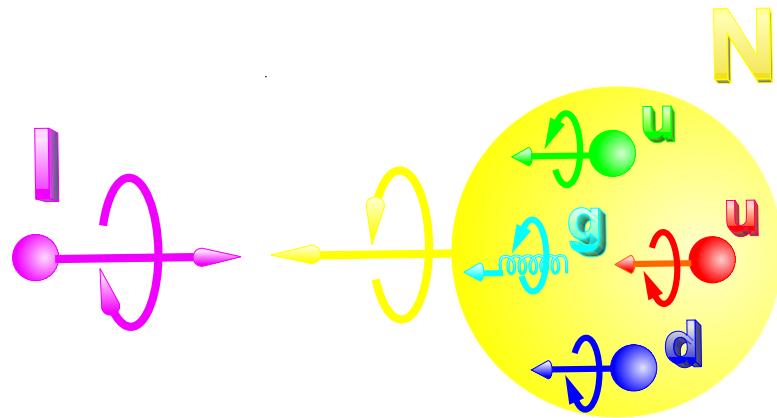
Dennison: resolves discrepancy between their results



concludes (June 16<sup>th</sup>, 1927) proton is a fermion of spin  $\frac{1}{2}$

Since then it is only getting much more complicated!

**quark model/QCD:** proton is a composite object

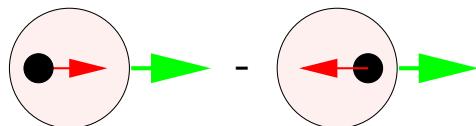


⇒ spin  $\frac{1}{2}$  must be the result of a subtle interplay of quark & gluon spins and their angular momentum

**quantitative description:** parton densities

parton model definition of “helicity densities”:

$$\Delta f(x) \equiv f_+^{N+}(x) - f_-^{N+}(x)$$

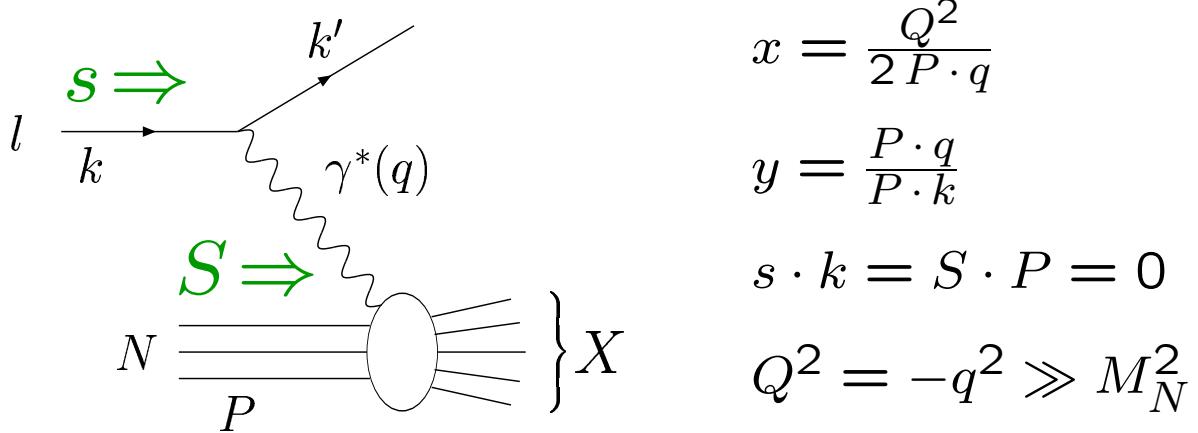


$f_{+(-)}^{N+}$ : “constituents” with same (opposite) helicity and momentum fraction  $x$  of the nucleon  $N$

**standard experiment:** “DIS- $\text{LN}$ -microscope”

$\tau_{\gamma^*} \ll$  time scale for parton-parton interactions

- lepton sees “snapshot” of nucleon
- lepton scatters off “free” quarks, incoherently



cross section:  $d^2\sigma/dxdy \propto \underbrace{\mathcal{L}_{\mu\nu}(k, q, s)}_{\text{leptonic}} \cdot \underbrace{\mathcal{W}^{\mu\nu}(P, q, S)}_{\text{hadronic}}$

$\mathcal{L}_{\mu\nu}(k, q, s)$  : fully calculable in QED

$\mathcal{W}^{\mu\nu}(P, q, S)$ : parametrizes our ignorance of hadron structure in terms of structure functions



$$g_1(x) = \frac{1}{2} \sum_q e_q^2 [\Delta q(x) + \Delta \bar{q}(x)]$$

for longitudinally polarized DIS

**Naive parton model picture is modified in full QCD:**

parton-parton interactions  $\Rightarrow$  spin & parton content of the nucleon depends on “resolving power” of probe

e.g. DIS: “resolution scale”  $R \simeq 1/\sqrt{Q^2}$

to explore scales  $R \lesssim \frac{1}{5} \text{ fm} \Leftrightarrow Q^2 \gtrsim 1 \div 2 \text{ GeV}^2$

## Main predictions of perturbative QCD

- applicable if hardness of probe is large ( $\alpha_s(Q^2) \ll 1$ )
- $\Delta f(x) \rightarrow \Delta f(x, Q^2)$ ; described by DGLAP eqs.:

$$\mu \frac{d}{d\mu} \begin{pmatrix} \Delta q(x, \mu) \\ \Delta g(x, \mu) \end{pmatrix} = \int_x^1 \frac{dz}{z} \begin{pmatrix} \Delta \mathcal{P}_{qq} & \Delta \mathcal{P}_{qg} \\ \Delta \mathcal{P}_{gq} & \Delta \mathcal{P}_{gg} \end{pmatrix}_{(z, \alpha_s(\mu))} \cdot \begin{pmatrix} \Delta q \\ \Delta g \end{pmatrix} \left( \frac{x}{z}, \mu \right)$$

$\Delta \mathcal{P}_{ij}$ : calculable spin-dep.  $j \rightarrow i$  “splitting” functions:

$$\Delta \mathcal{P}_{ij}(z, \alpha_s) = \frac{\alpha_s}{2\pi} \Delta P_{ij}^{(0)}(z) + \left( \frac{\alpha_s}{2\pi} \right)^2 \Delta P_{ij}^{(1)}(z) + \dots$$

Ahmed, Ross  
Altarelli, Parisi              Mertig, van Neerven  
                                    Vogelsang

- scaling violations in DIS:  $g_1(x) \rightarrow g_1(x, Q^2)$

$$g_1(x, Q^2) = \frac{1}{2} \sum_q e_q^2 \left\{ \left[ \Delta q + \Delta \bar{q} \right] (x, Q^2) \right.$$

$$+ \left| \begin{array}{c} \text{wavy line} \\ \diagdown \quad \diagup \\ \text{horizontal line} \end{array} \right. + \dots \left| \begin{array}{c} ^2 \\ \otimes \left[ \Delta q + \Delta \bar{q} \right] (x, Q^2) + \dots \end{array} \right. \right.$$

$$+ \left| \begin{array}{c} \text{wavy line} \\ \diagup \quad \diagdown \\ \text{horizontal line} \end{array} \right. + \dots \left| \begin{array}{c} ^2 \\ \otimes \Delta g(x, Q^2) + \dots \end{array} \right. \right\} + \mathcal{O}\left(\frac{1}{Q^n}\right)$$

$\mathcal{O}(\alpha_s)$  QCD corrections (“coefficient functions”  $\Delta C_{q,g}$ )

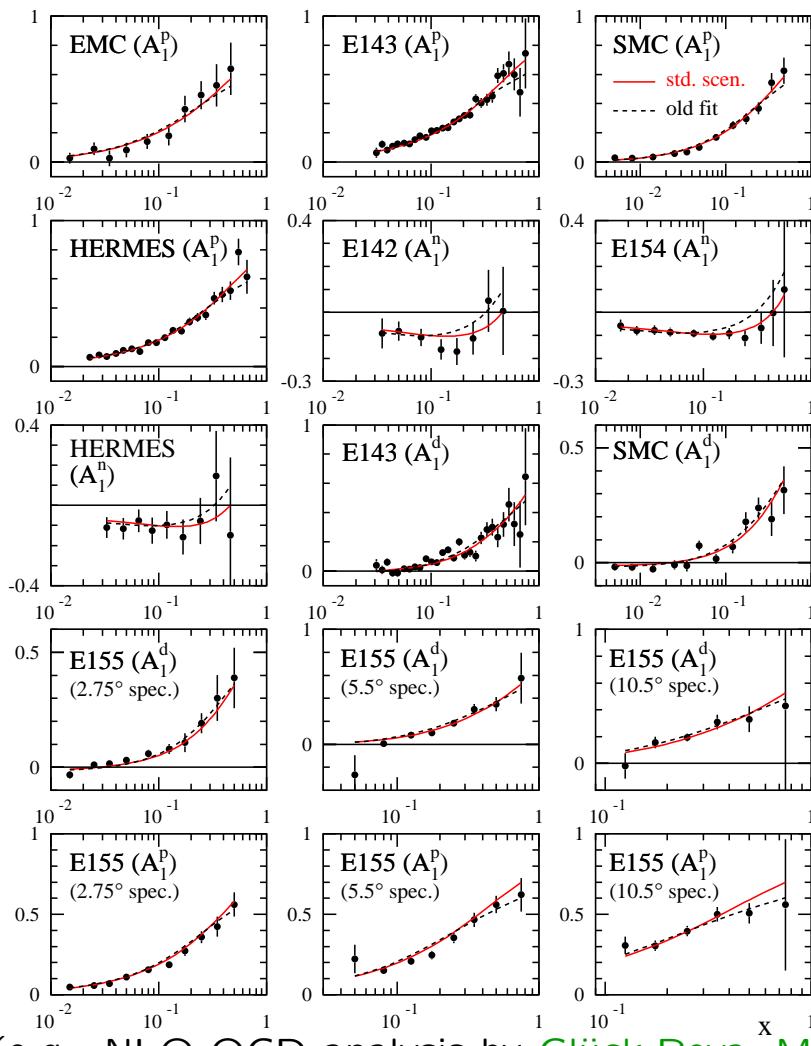
Kodaira; Ratcliffe; Bodwin, Qiu; van Neerven, Zijlstra; ...

More than 20 years of beautiful data on polarized DIS



## What have we learned so far?

- NLO QCD fits work extremely well:



(e.g., NLO QCD analysis by Glück, Reya, MS, Vogelsang)

- nucleon spin-1/2 “sum rule”:

Jaffe, Manohar; Ji

$$\frac{1}{2} = \underbrace{\frac{1}{2} \sum_q \int_0^1 [\Delta q + \Delta \bar{q}] (x, Q^2) dx}_{\text{striking finding : } \approx 0.1} + \underbrace{L_q(Q^2) + J_g(Q^2)}_{\text{unknown}}$$

contrary to naive quark model

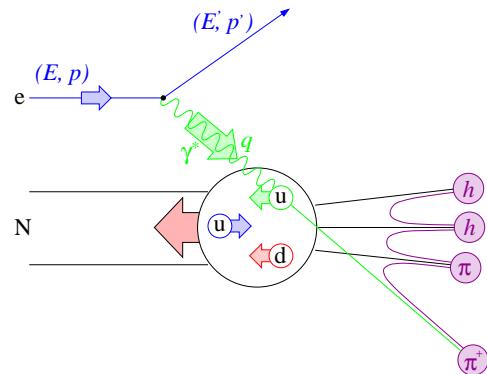
$L_q$ : quark *orbital* angular momentum

$J_g$ : gluon *total* angular momentum, contains  $\int_0^1 \Delta g(x, Q^2) dx$

- recent exp. progress in semi-incl. DIS:

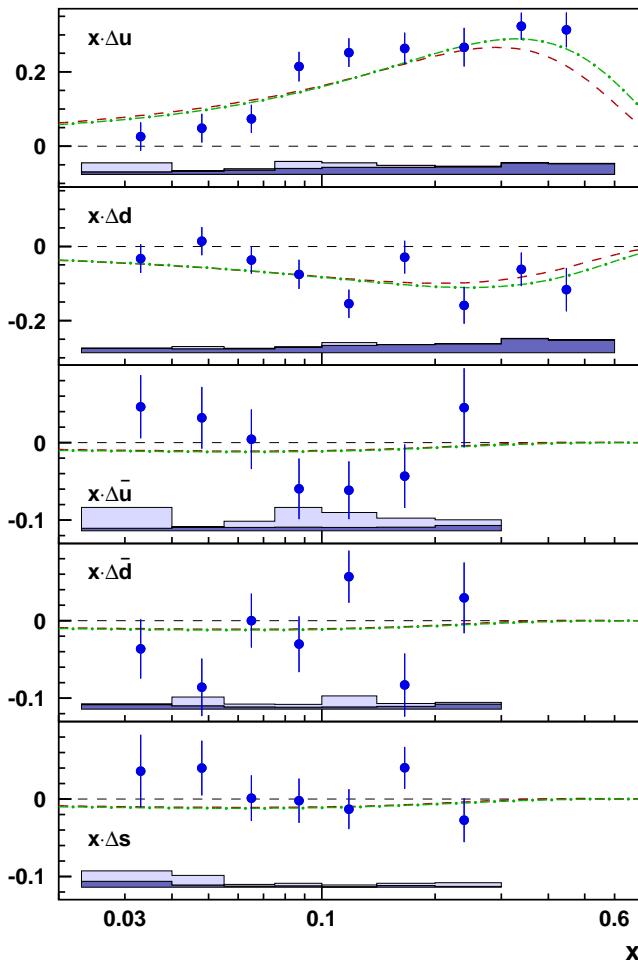


**idea:** flavor content of the observed hadron  $\leftrightarrow$  flavor of the struck quark via frag. fcts.  $D_q^h(z)$   
 $\Rightarrow q/\bar{q}$  & flavor separation possible



LO spin asymmetry:  $(Q^2 \text{ dependence suppressed})$

$$A_1^h(x, z) = \frac{g_1^h(x, z)}{F_1^h(x, z)} = \frac{\sum_q e_q^2 \left[ \Delta q(x) D_q^h(z) + \Delta \bar{q}(x) D_{\bar{q}}^h(z) \right]}{\sum_q e_q^2 \left[ q(x) D_q^h(z) + \bar{q}(x) D_{\bar{q}}^h(z) \right]}$$



**considerable interest:**

SU(2) breaking in sea  
 Diakonov et al.; Fries et al.;  
 Kumano; Thomas et al.; ...

strange polarization  $\Delta_s$   
 Brodsky et al.; Ellis et al.;  
 ...

**points of discussion:**

positive  $\Delta_s$ ?  
 Bass; Leader, Stamenov

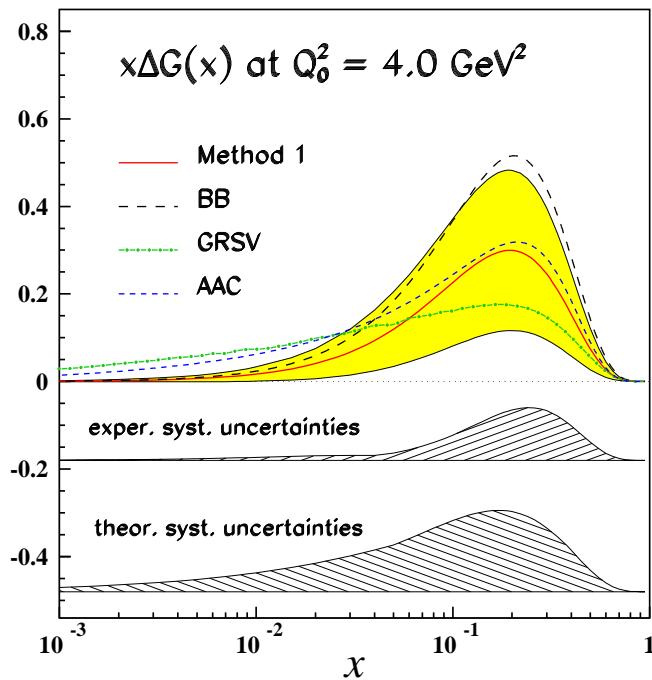
use of  $D_q^h(z)$  at low  $Q^2$ ?  
 Kretzer et al.; Kotzinian

higher twist?

But we are still not out of the woods -

the jury is still out on some **crucial questions**:

- need more information on  $\Delta q$ ,  $\Delta \bar{q}$  & flavor separation
- present data cover only small range in  $x$  and  $Q^2$ 
  - DIS constrains  $\Delta g$  mainly through scaling violations  
→  $\Delta g(x, Q^2)$  completely/largely unknown



taken from a recent fit by Böttcher, De Nardo

- impossible to really verify predicted  $\ln Q^2$  behavior
- still don't know very much about spin- $\frac{1}{2}$  "sum rule"
- haven't tested universality of polarized parton densities
- even worse: transverse polarization



 - Spin Program

## II. Spin Physics at



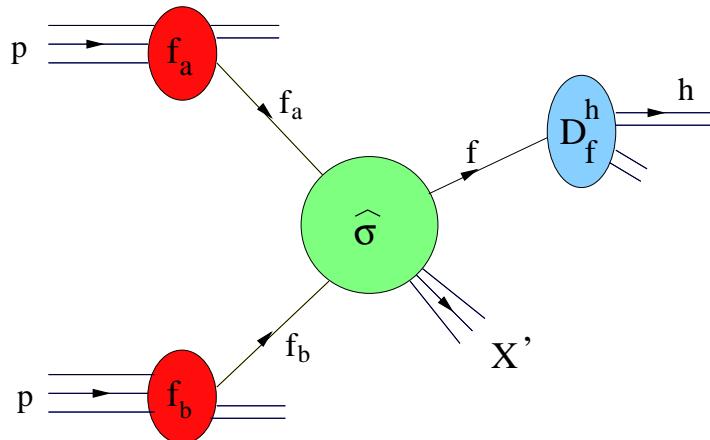
**starting point:** universality of parton densities

- foundation of predictive power of perturbative QCD
- makes notion of “nucleon structure” meaningful

**factorization @ work:** high- $p_T$  processes

Libby, Sterman; Ellis et al.; Amati et al.; Collins et al.; ...

e.g., high- $p_T$  single-inclusive hadron production



**long-distance**

from exp.;  $\mu$ -dep.:  $d\sigma/d\mu = 0$  (pQCD)



$$\frac{d\sigma^{pp \rightarrow \pi X}}{d\mathcal{P}} = \sum_{abc} \int dx_a dx_b dz_c f_a(x_a, \mu_f) f_b(x_b, \mu_f) D_c^\pi(z_c, \mu_{f'})$$

$$\times \frac{d\hat{\sigma}^{ab \rightarrow cX'}}{d\mathcal{P}}(x_a P_a, x_b P_b, P^\pi/z_c, \mu_f, \mu_{f'}, \mu_r) + \mathcal{O}\left(\frac{\lambda}{p_T}\right)^n$$

**short-distance**

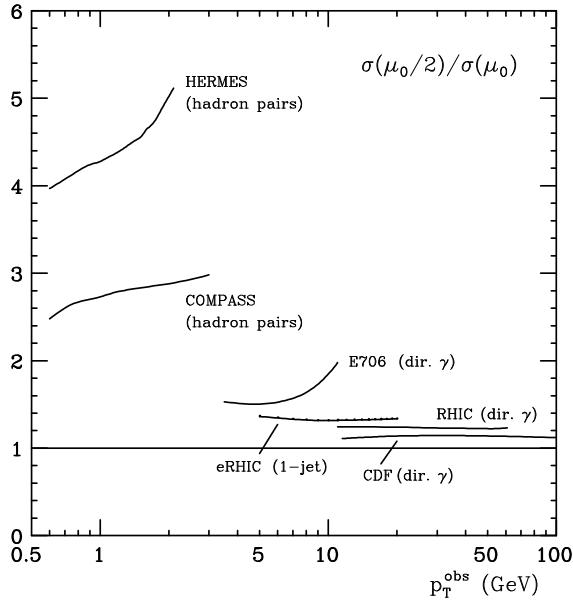
calculable in pQCD: power series in  $\alpha_s$

power corr.  
neglected

- $\mathcal{P}$ : appropriate set of kin. variables ( $p_T, y, \dots$ )
- arbitrary scales  $\mu_{f,f',r}$ : separate long- and short-dist. physics

## QCD in a collider environment: many advantages ...

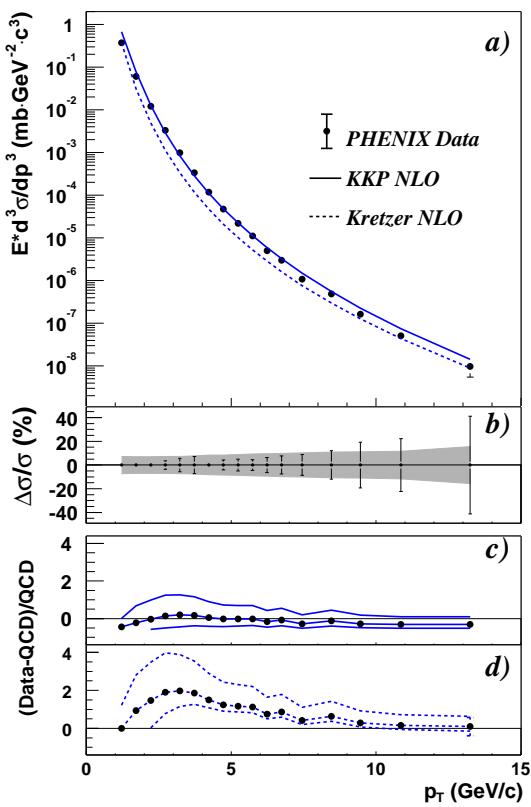
- large c.m.s. energy → high  $p_T$  accessible
  - domain of perturbative QCD:  $\alpha_s(\mu_r \simeq p_T) \ll 1$
  - scale dependence ↔ reliability of pQCD calc.:



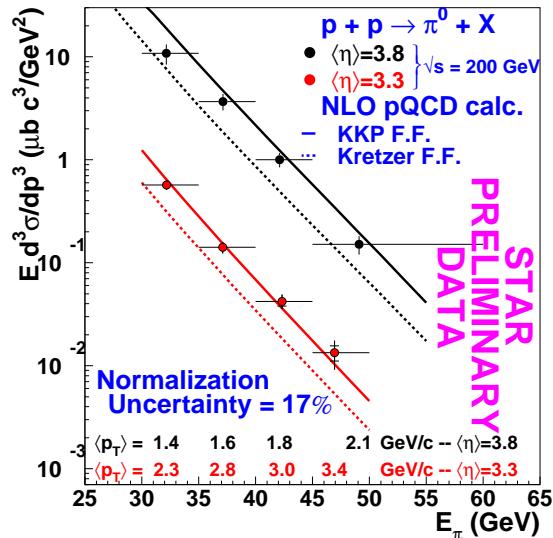
ratios of “x-secs” for  
two choices of  $\mu_f$

← should be  $\simeq 1$

- a success story in describing *unpolarized* collisions

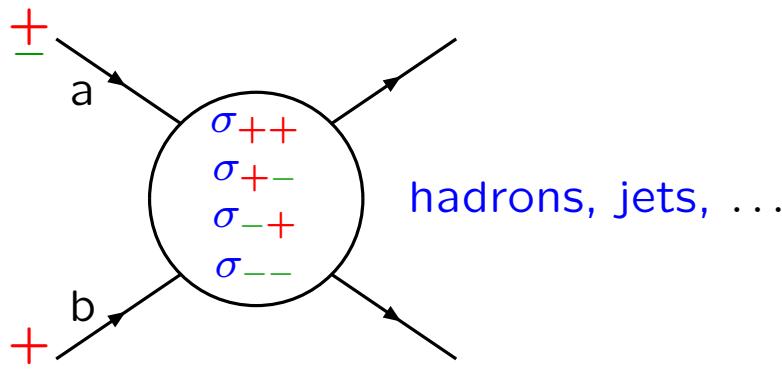


1<sup>st</sup> examples from RHIC:  
high  $p_T$  pion production



agrees well with NLO pQCD

# Longitudinal Polarization



$$\rightarrow \frac{1}{2} (\sigma_{++} + \sigma_{+-}) \equiv \sigma \quad (\text{usual unpol. cross section})$$

$$\frac{1}{2} (\sigma_{++} - \sigma_{+-}) \equiv \Delta\sigma = \Delta f_a \otimes \Delta f_b \otimes \Delta\hat{\sigma}$$

experimentally relevant:

$$A_{LL} \equiv \frac{d\Delta\sigma/dP}{d\sigma/dP}$$

## ultimate goal:

pin down all aspects of helicity pdfs, in particular  $\Delta g$

reaction	LO subprocesses	partons probed	$x$ -range
$pp \rightarrow \text{jets } X$	$q\bar{q}, qq, qg, gg \rightarrow \text{jet } X$	$\Delta q, \Delta g$	$x \gtrsim 0.03$
$pp \rightarrow \pi X$	$q\bar{q}, qq, qg, gg \rightarrow \pi X$	$\Delta q, \Delta g$	$x \gtrsim 0.03$
$pp \rightarrow \gamma X$	$qg \rightarrow q\gamma, q\bar{q} \rightarrow g\gamma$	$\Delta g$	$x \gtrsim 0.03$
$pp \rightarrow Q\bar{Q}X$	$gg \rightarrow Q\bar{Q}, q\bar{q} \rightarrow Q\bar{Q}$	$\Delta g$	$x \gtrsim 0.01$
$pp \rightarrow W^\pm X$	$q\bar{q}' \rightarrow W^\pm$	$\Delta u, \Delta \bar{u}, \Delta d, \Delta \bar{d}$	$x \gtrsim 0.06$



- test of fundamental QCD spin interactions
- verify universality of pdfs for the 1<sup>st</sup> time
- test predicted  $\ln Q^2$  dependence of pdfs for the 1<sup>st</sup> time
- perhaps learn more about spin-1/2 “sum rule”

# Complication: NLO QCD Corrections

## Why NLO?

- expect reduced dependence on unphysical scales  $\mu_{f,f,r'}$
- NLO corrections often sizable
- more reliable ang./ $p_T$  distributions, jet definition, etc.
- genuine NLO processes may dilute sensitivity to  $\Delta g$

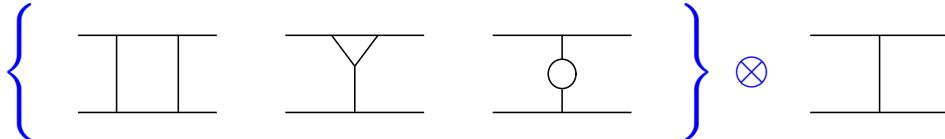
## What are NLO QCD corrections all about?

at  $\mathcal{O}(\alpha_s^2)$  (LO) one has:

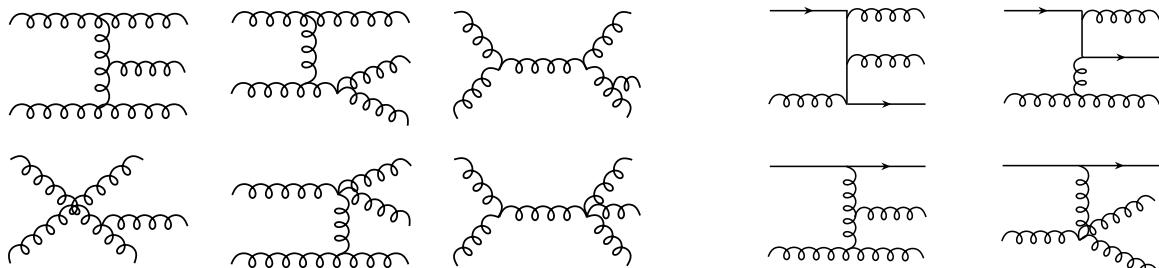
relevant  $2 \rightarrow 2$   parton-parton scattering processes

at  $\mathcal{O}(\alpha_s^3)$  (NLO) one has:

- interference of 1-loop and LO Born amplitudes:



- $2 \rightarrow 3$   parton-parton scattering processes, e.g.:



$gg \rightarrow ggg$

$qg \rightarrow qgg$

UV poles (loops): removed by *renormalization* → scale  $\mu_r$

IR poles: cancel in sum of virtual and  $2 \rightarrow 3$  diagrams

collinear poles: removed by *factorization* → scales  $\mu_{f,f'}$

**good news:** impressive progress recently

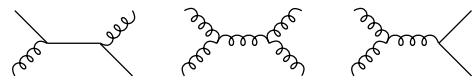
<b>evol. kernels</b>	$\Delta P_{ij}$ (spacelike) $\Delta P_{i\gamma}$ (photon) $\Delta P_{ij}$ (timelike)	NLO NLO NLO		Mertig, van Neerven; Vogelsang Vogelsang, MS Vogelsang, MS
<b>DIS</b>	$\vec{e}\vec{p} \rightarrow e'X$	NLO NNLO		Kodaira et al.; Ratcliffe; Bodwin, Qiu van Neerven, Zijlstra
<b>SIDIS</b>	$\vec{e}\vec{p} \rightarrow e'H + X$ $\vec{e}\vec{p} \rightarrow e'\vec{H} + X$ $\vec{e}\vec{p} \rightarrow e'\vec{H} + X$	NLO NLO NLO		Chiappetta et al.; De Florian et al. De Florian et al. De Florian et al.
<b>hadrons</b>	$\vec{p}\vec{p} \rightarrow H + X$ $\vec{\gamma}\vec{p} \rightarrow H + X$ $\vec{p}\vec{p} \rightarrow \vec{H} + X$	NLO NLO NLO	•	De Florian; Jäger et al. Jäger, MS, Vogelsang Jäger et al. (soon!)
<b>jets</b>	$\vec{p}\vec{p} \rightarrow \text{jet(s)} + X$ $\vec{\gamma}\vec{p} \rightarrow \text{jet(s)} + X$ $\vec{e}\vec{p} \rightarrow \text{jet(s)} + X$	NLO NLO NLO	•	De Florian et al.; Jäger et al. (soon!) De Florian, Frixione Mirkes, Willfahrt
<b>prompt <math>\gamma</math></b>	$\vec{p}\vec{p} \rightarrow \gamma + X$ $\vec{p}\vec{p} \rightarrow \gamma\gamma + X$ $\vec{e}\vec{p} \rightarrow \gamma + X$ $\vec{p}\vec{p} \rightarrow \vec{\gamma} + X$	NLO NLO NLO NLO	•	Gordon, Vogelsang; Contogouris et al. Coriano, Gordon Contogouris et al. Vogelsang
<b><math>\gamma + \text{jet}</math></b>	$\vec{p}\vec{p} \rightarrow \gamma + \text{jet} + X$ $\vec{\gamma}\vec{p} \rightarrow \gamma + \text{jet} + X$	NLO NLO		Gordon Gordon
<b><math>\gamma + \text{charm}</math></b>	$\vec{p}\vec{p} \rightarrow \gamma + c + X$	NLO		Berger et al. ( $m_c = 0$ )
<b>heavy quarks</b>	$\vec{p}\vec{p} \rightarrow Q\bar{Q}X$ $\vec{\gamma}\vec{p} \rightarrow Q\bar{Q}X$	NLO NLO	•	Bojak, MS Bojak, MS; Contogouris et al.
<b>Drell-Yan</b>	$\vec{p}\vec{p} \rightarrow (\gamma^*)X$	NLO NNLO		Weber; Gehrmann; Smith et al.
<b>vector bosons</b>	$\vec{p}\vec{p} \rightarrow (Z^0, W^\pm)X$ $\vec{p}\vec{p} \rightarrow (Z^0, W^\pm)X$	NLO NLO	•	Weber; Gehrmann Weber; Gehrmann



details/results: → next slides

## Processes one-by-one

- **high- $p_T$  hadrons:**

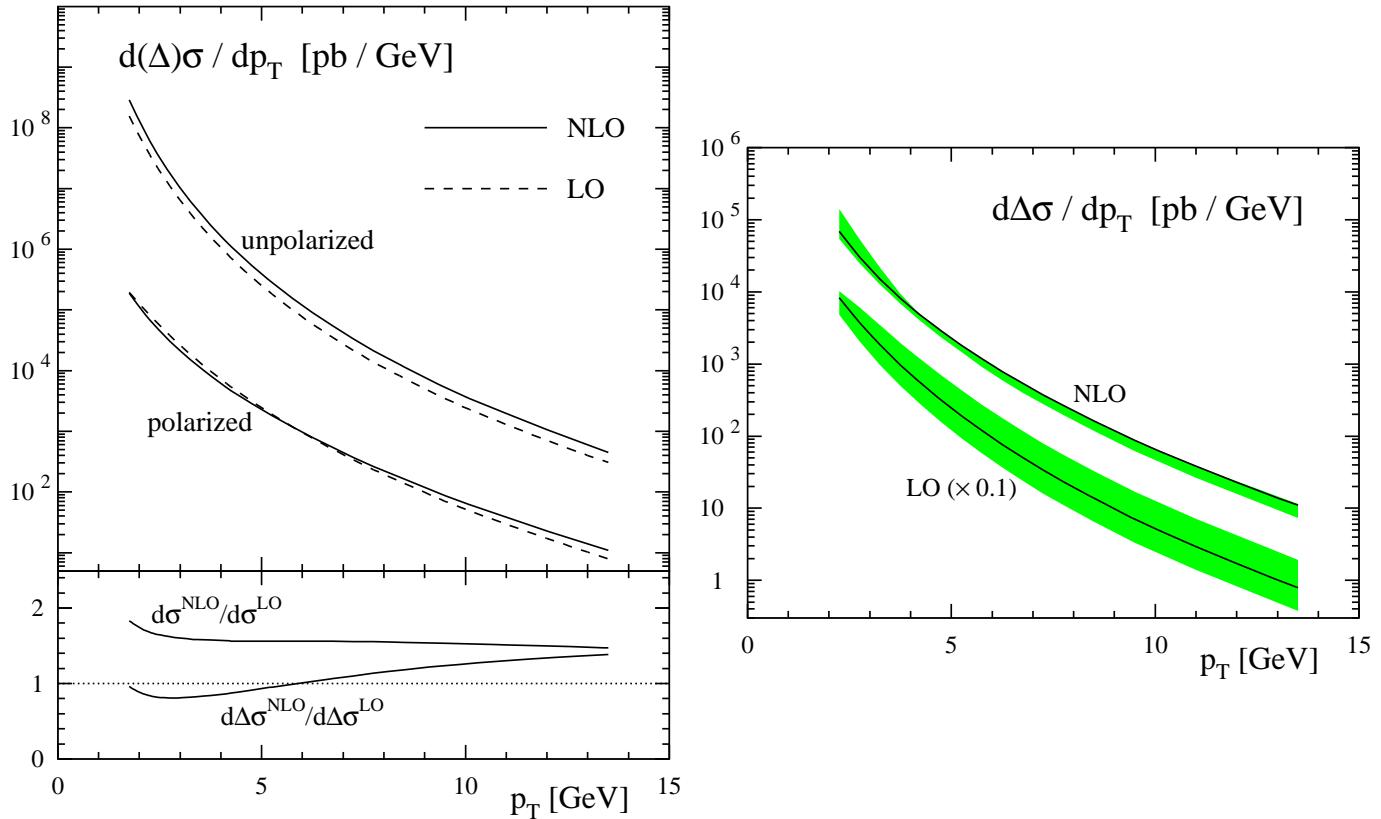


NLO: Jäger, MS, Schäfer, Vogelsang; de Florian

- recall:**
- unpol. measurements agree well with pQCD
  - $D^\pi(z)$  set of Kniehl, Kramer, Pötter favored
  - exciting prospects for polarized measurement

- results:**
- $d(\Delta)\sigma^{\text{NLO}}/d(\Delta)\sigma^{\text{LO}}$  reasonably small
  - theor. uncertainties much reduced in NLO

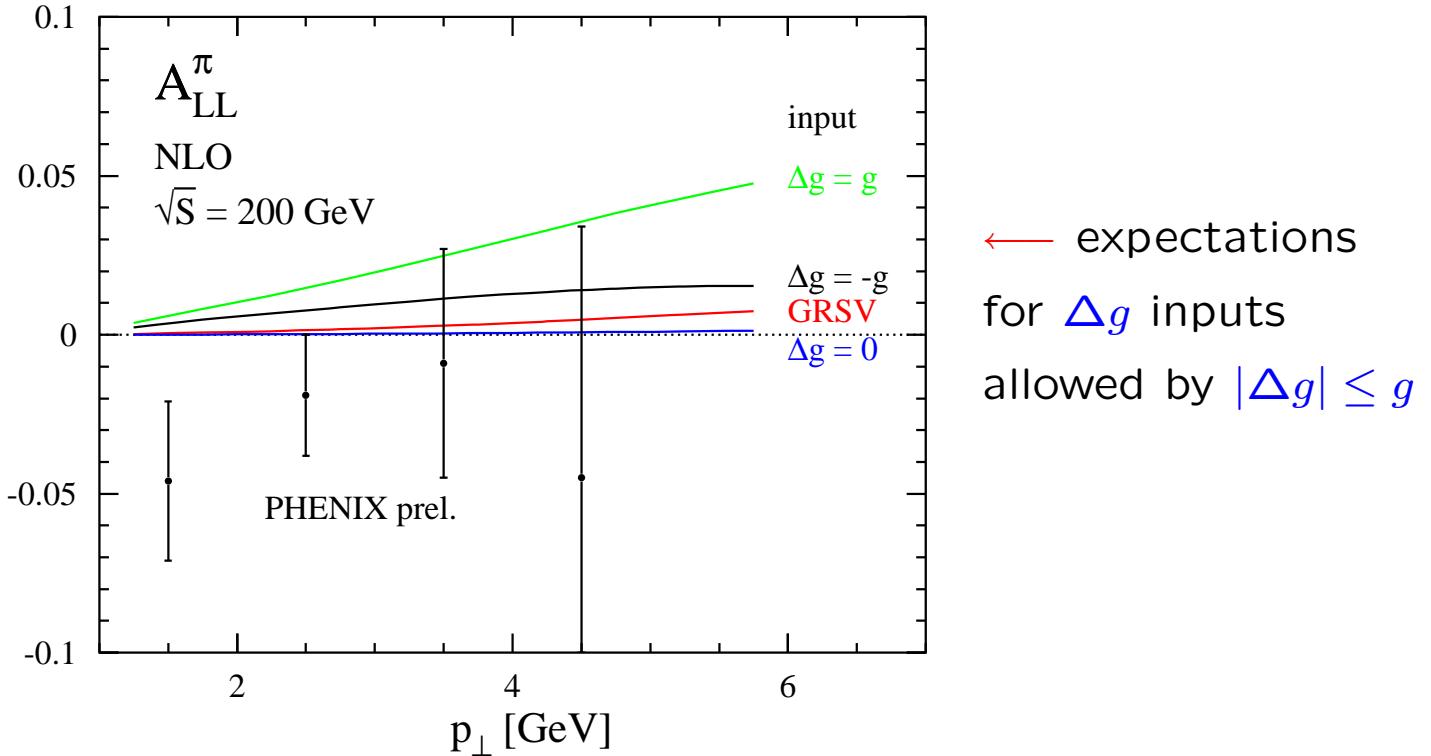
$\pi^0$ -production,  $\sqrt{S} = 200 \text{ GeV}$ ,  $|\eta| \leq 0.38$



[figs. taken from Jäger et al.]

# exciting development: first prel. results by

[pol. scaling error ( $\sim 30\%$ ) not included]



## naive analysis:

partonic spin asymmetries:

$$\begin{aligned} gg \rightarrow gg & \quad \hat{a}_{LL} > 0 \\ gg \rightarrow q\bar{q} & \quad \hat{a}_{LL} - 1 \\ gq \rightarrow gq & \quad \hat{a}_{LL} > 0 \end{aligned}$$

$$\langle z \rangle = 0.3 \div 0.4 \text{ where } D_q^\pi > D_g^\pi$$

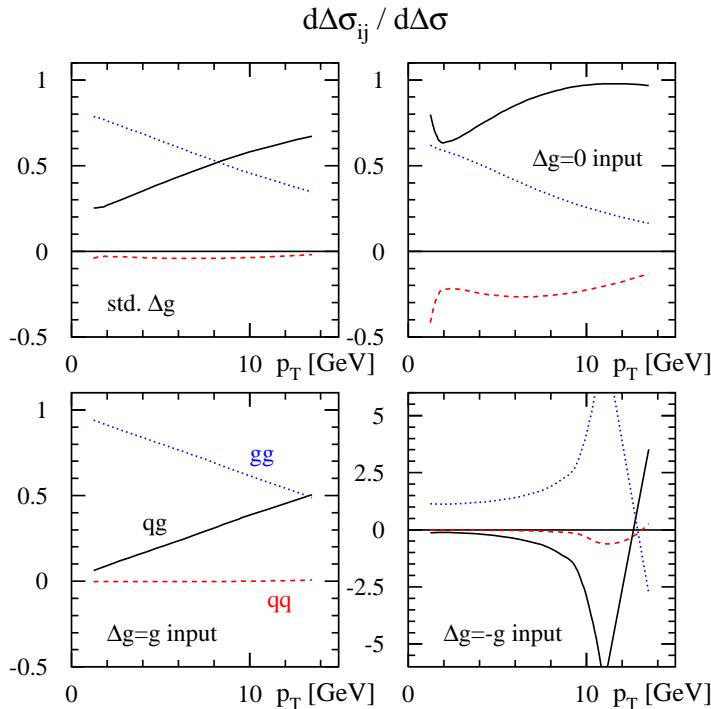
→ conclude:  $gg \rightarrow q\bar{q}$  resp. for neg.  $A_{LL}^\pi$

**but**

$$\Delta \hat{\sigma}_{gg \rightarrow gg} \simeq 160 \Delta \hat{\sigma}_{gg \rightarrow q\bar{q}} !! \text{ (for } \eta \simeq 0)$$

# a closer look: can $A_{LL}^\pi$ be negative?

Jäger, Kretzer, MS, Vogelsang



**subprocess contrib.:**

$gg$  process dominates  
for  $p_T \lesssim 10$  GeV

$qg$  process takes over  
for  $p_T \gtrsim 10$  GeV

$qq$  channel always small

take  $x_T^2 \equiv 4p_T^2/S$  moments of  $\eta$ -integrated  $d\Delta\sigma$ :

$$\Delta\sigma^\pi(N) = (\Delta g^{N+1})^2 \mathcal{A}^N + 2\Delta g^{N+1} \mathcal{B}^N + \mathcal{C}^N$$

$\uparrow$                      $\uparrow$                      $\uparrow$   
 $gg$                      $qg$                      $qq$

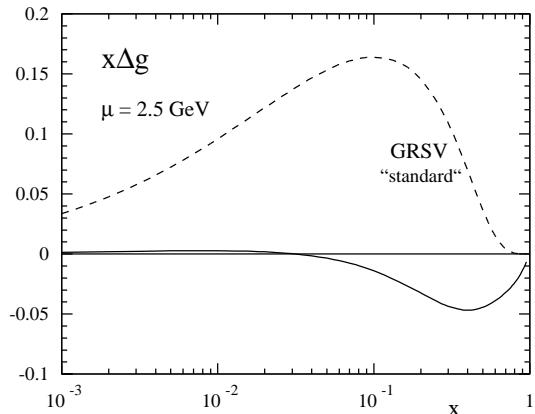
$$\text{minimize : } \Delta\sigma^\pi(N) \Big|_{\min} = -\frac{(\mathcal{B}^N)^2}{\mathcal{A}^N} + \mathcal{C}^N$$

→ negative, but tiny lower bound  $A_{LL}^\pi \Big|_{\min} \simeq \mathcal{O}(-10^{-3})$

**resulting  $\Delta g$ :**

node! →  $\Delta g(x_a)\Delta g(x_b) < 0$

confirmed by full “global analysis”



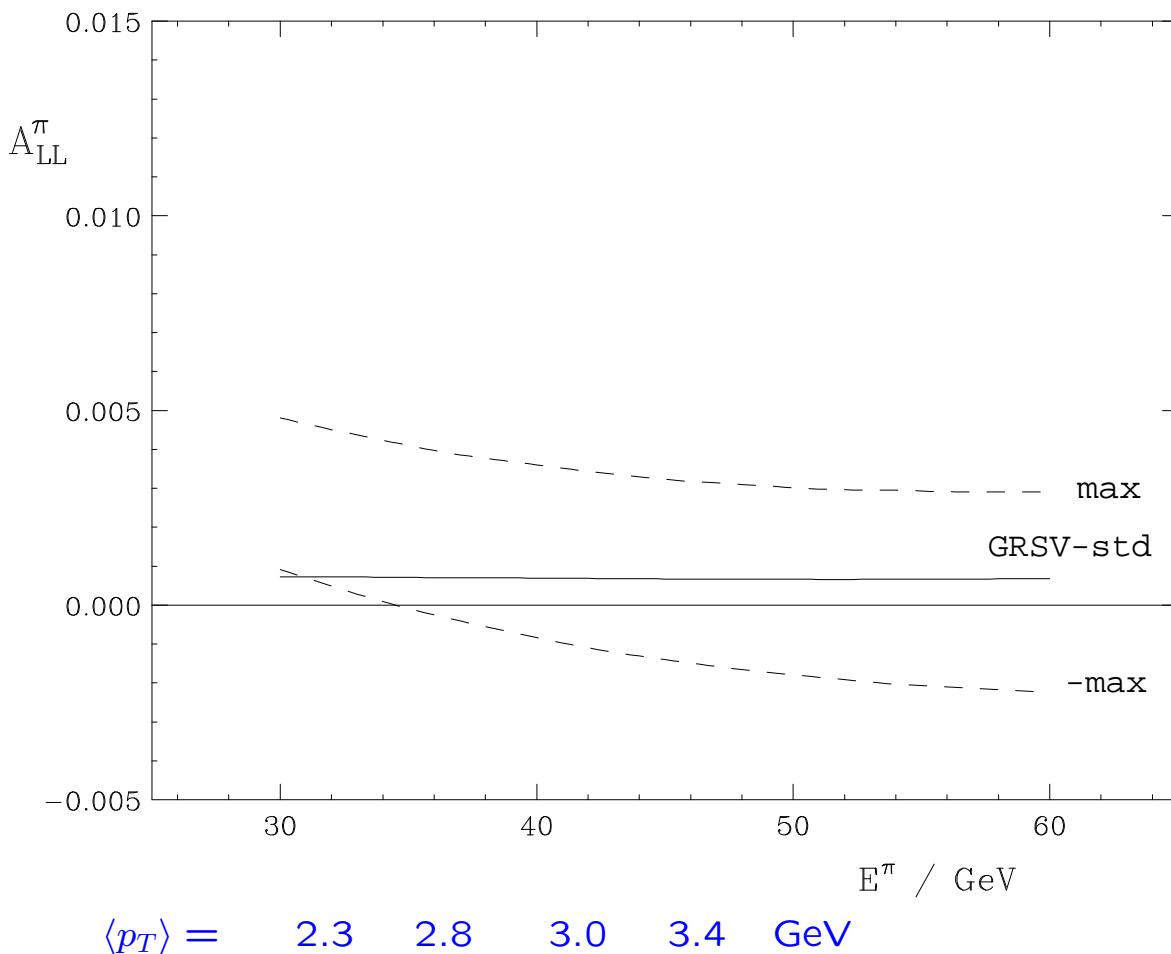
**recall:** successful measurement of unpolarized cross section in forward region  
agrees with pQCD  
→  $A_{LL}^\pi$  measurement at large rapidities?

**kinematics:**  $\langle\eta\rangle \simeq 3.5$ ;  $\langle p_T \rangle \simeq 1.5 \div 3$  GeV;  $\langle z \rangle \simeq 0.7 \div 0.8$

**expect:** dominance of  $qg$  subprocess sets in  
→ sign/size of  $A_{LL}^\pi$  tied to sign/size of  $\Delta g$

**results:**

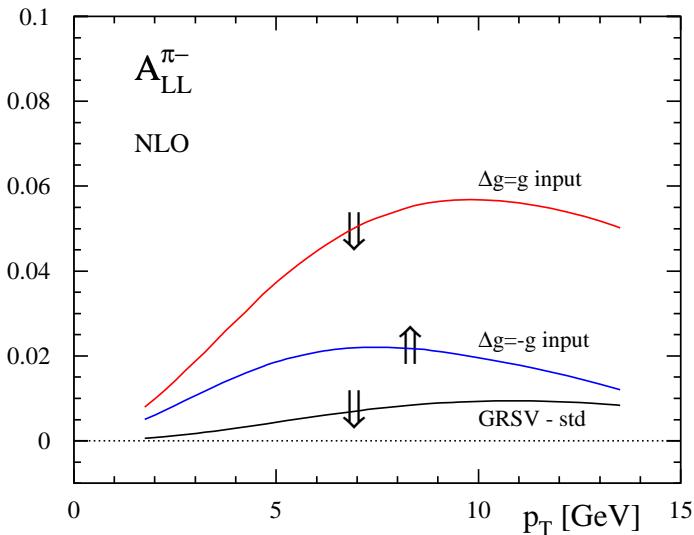
$$\sqrt{S} = 200 \text{ GeV}, \eta = 3.3$$



# long-term prospects on $A_{LL}^{\pi}$ for :

Jäger, MS, Vogelsang; Saito

compare  $A_{LL}^{\pi^{\pm}}$  to  $A_{LL}^{\pi^0}$



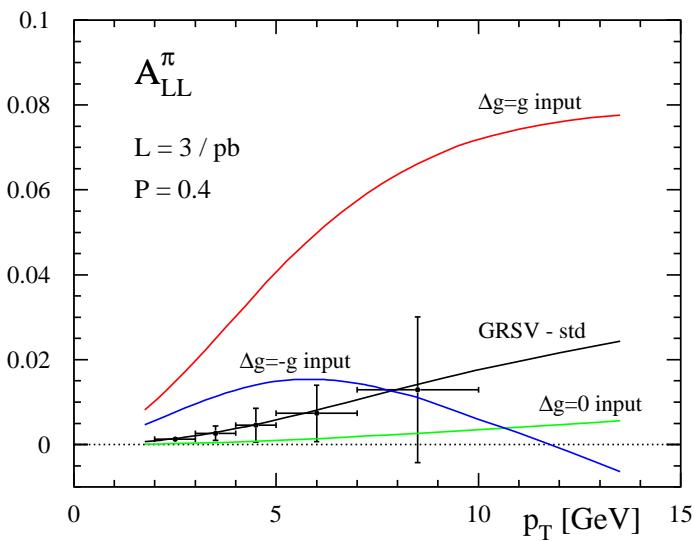
$\pi^-$

**idea:**

for  $p_T \gtrsim 5$  GeV:

$qg$  starts to dominate

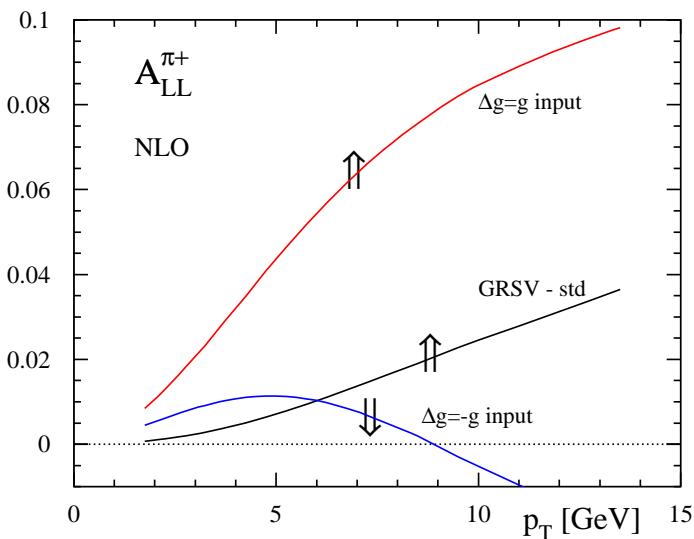
and  $D_u^{\pi^+} > D_u^{\pi^-}$ ,  $D_g^{\pi^+} = D_g^{\pi^-}$



**expect:**

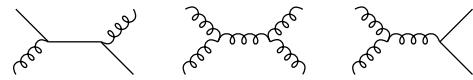
sensitivity to sign of  $\Delta g$

e.g. pos.  $\Delta g$ :  $A_{LL}^{\pi^+} > A_{LL}^{\pi^-}$



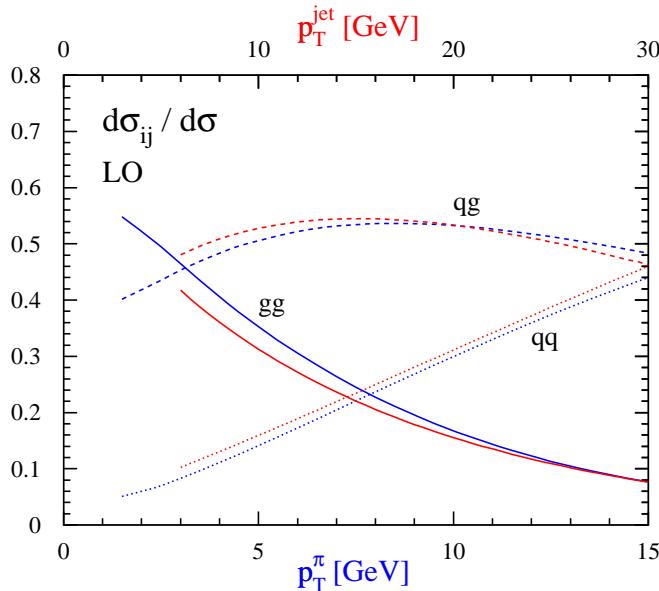
$\pi^+$

- **high- $p_T$  jets:**



NLO: de Florian, Frixione, Signer, Vogelsang; Jäger, MS, Vogelsang

proceeds through the same subprocesses as  $\pi$ -production



$\pi$ 's have roughly  $\langle z \rangle \simeq 0.5$ :

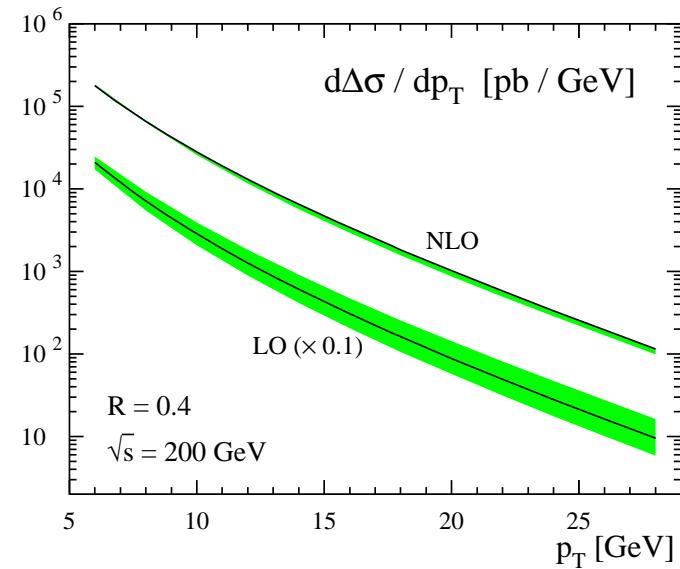
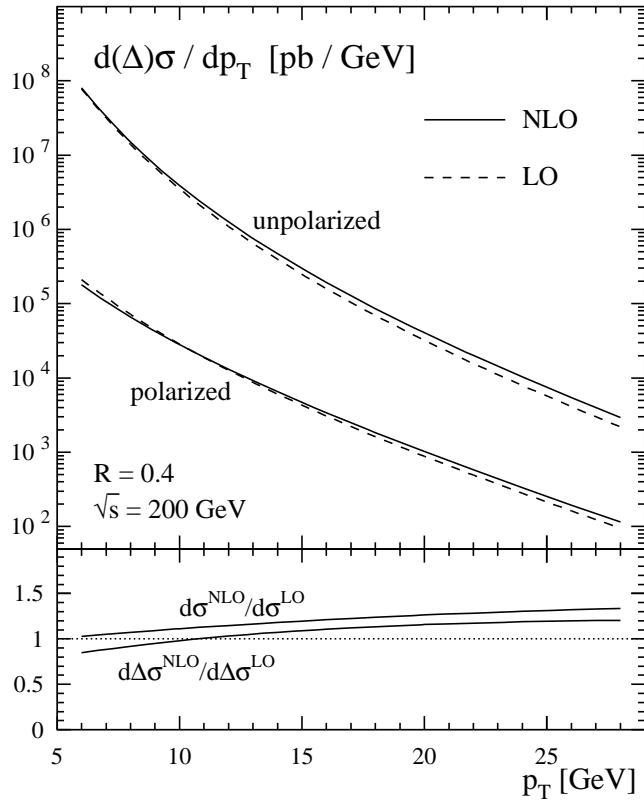
$\pi$  with  $p_T \simeq$  jet with  $2p_T$   
 $\rightarrow$  jet and  $\pi$  results similar

### advantages of jets:

- much higher rates
- no uncertainties from  $D(z)$

**results:** theor. uncertainties even smaller than for hadrons

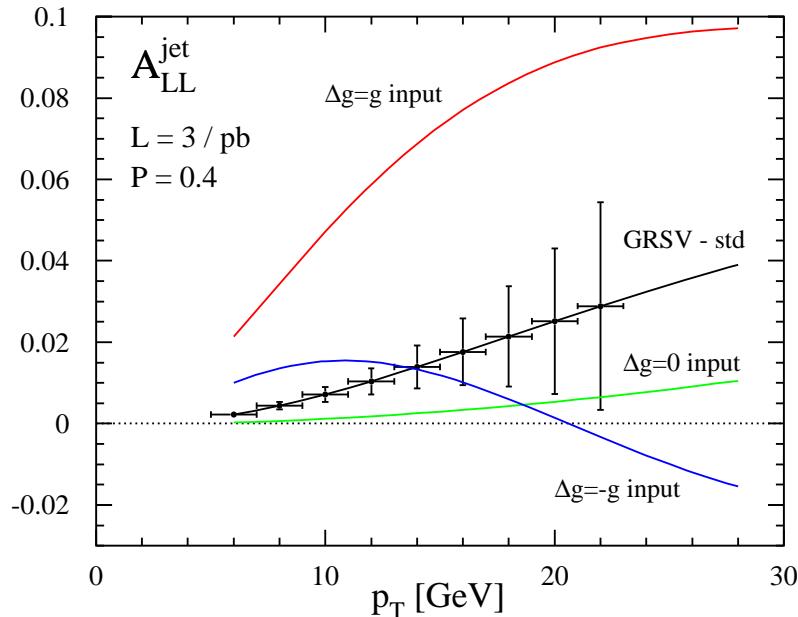
$$\sqrt{S} = 200 \text{ GeV}, R_{\text{cone}} = 0.4 \text{ (SCA)}, |\eta| \leq 1$$



[figs. taken from Jäger, MS, Vogelsang]

# expectations for :

$\sqrt{S} = 200 \text{ GeV}$ ,  $R_{\text{cone}} = 0.4$  (SCA),  $0 \leq \eta \leq 1$

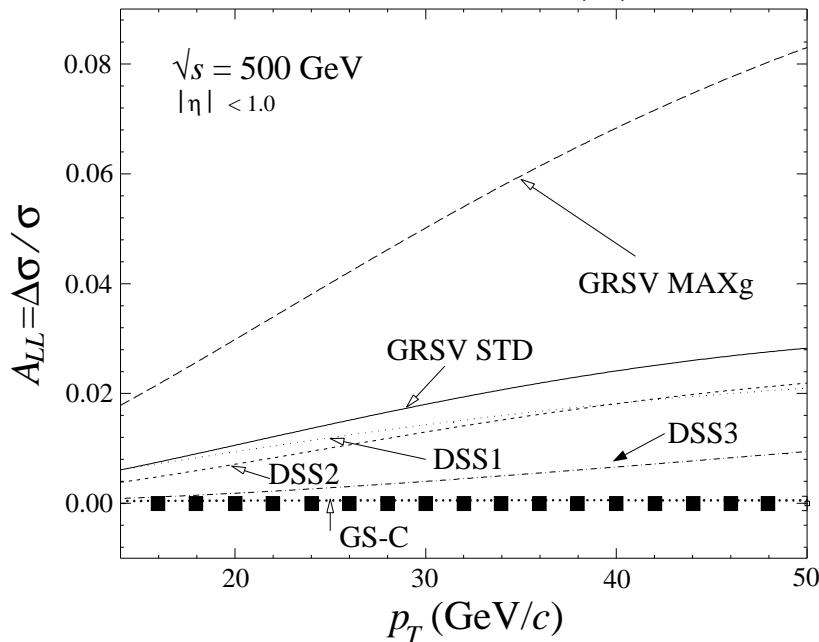


[fig. taken from Jäger, MS, Vogelsang]

estimate of statistical errors:  $\delta A \simeq \frac{1}{P_1 P_2} \times \frac{1}{\sqrt{\mathcal{L} \sigma_{\text{bin}}}}$   
 [with  $P_1 = P_2 = 0.4$  (beam pol.) and  $\mathcal{L} = 3 \text{ pb}^{-1}$ ]

and in the long run:

$\sqrt{S} = 500 \text{ GeV}$ ,  $D = 1$  (ES),  $|\eta| \leq 1$ ,  $\mathcal{L} = 800 \text{ pb}^{-1}$



[fig. taken from de Florian, Frixione, Signer, Vogelsang]

# summary on inclusive jets and hadrons:

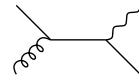
- excellent near-term prospects for a first determination of  $\Delta g$
- important: unpolarized “benchmark x-secs” so far in perfect agreement with pQCD [even down to unexpectedly small  $p_T$ ]
- first exciting results from  if large & negative  $A_{LL}$  persists: new spin surprise; effect due to intr.  $k_T$  or ? important to collect more data **soon**

## long-term goals:

need much higher luminosities and/or  $\sqrt{S} = 500 \text{ GeV}$

- prompt photons: cross-check of  $\Delta g$  results
- heavy flavors: extend  $\Delta g$  to smaller  $x$
- $W^\pm$  bosons: further information on  $\Delta q, \Delta \bar{q}$

## ● prompt photons (plus jets):



NLO: Gordon, Vogelsang; Contogouris et al.; Frixione, Vogelsang

**idea:**  $qg \rightarrow q\gamma$  dominant → clean signal for  $\Delta g$

**results:** • some trouble with unpol. fixed target data  
but collider data O.K.

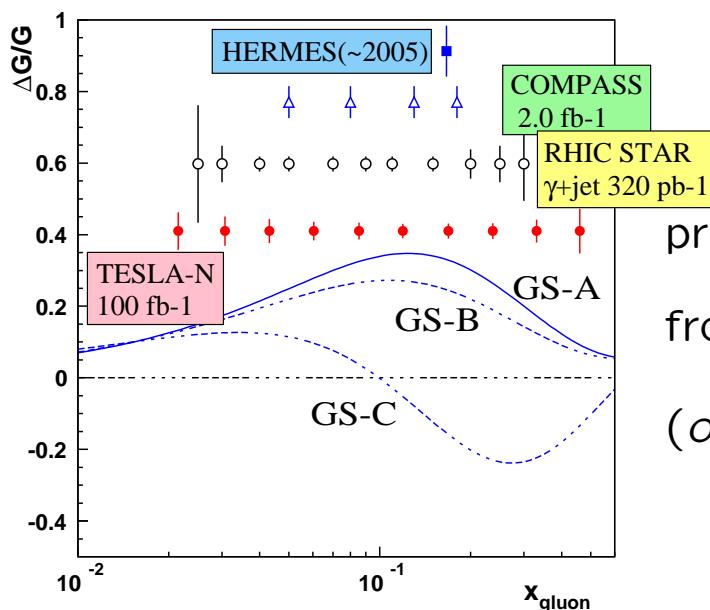
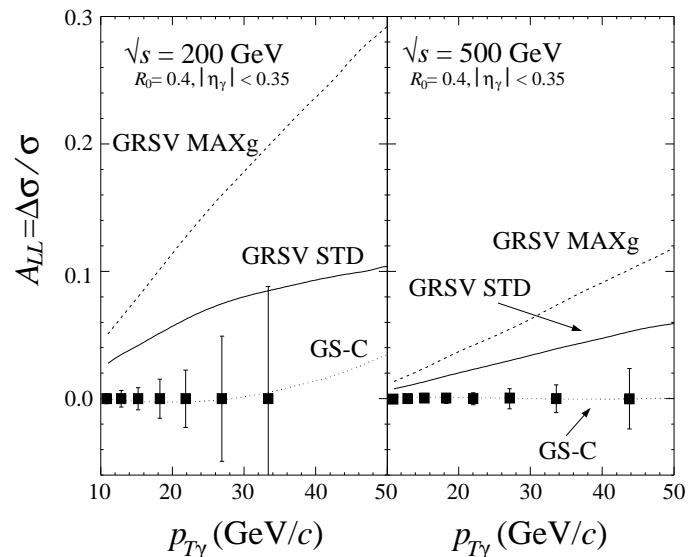
- theoretical efforts to go beyond NLO  
[soft gluon/threshold (and  $k_T$ ) resummations]  
Laenen et al.; Catani et al.; Li; Kidonakis Owens; ...

$A_{LL}^\gamma(p_T)$  for **PHENIX**:  
Frixione, Vogelsang

→ expected sensitivity to  $\Delta g$

not shown:

- reduced scale dependence  
in NLO ✓



projected accuracy for  $\Delta g/g$   
from  $\gamma + \text{jet}$  at **STAR**  
(only for 200 GeV meas.)

- **heavy flavor production:**

NLO: Bojak, MS

**idea:** sensitivity to  $\Delta g$  through dominance of  $gg \rightarrow Q\bar{Q}$

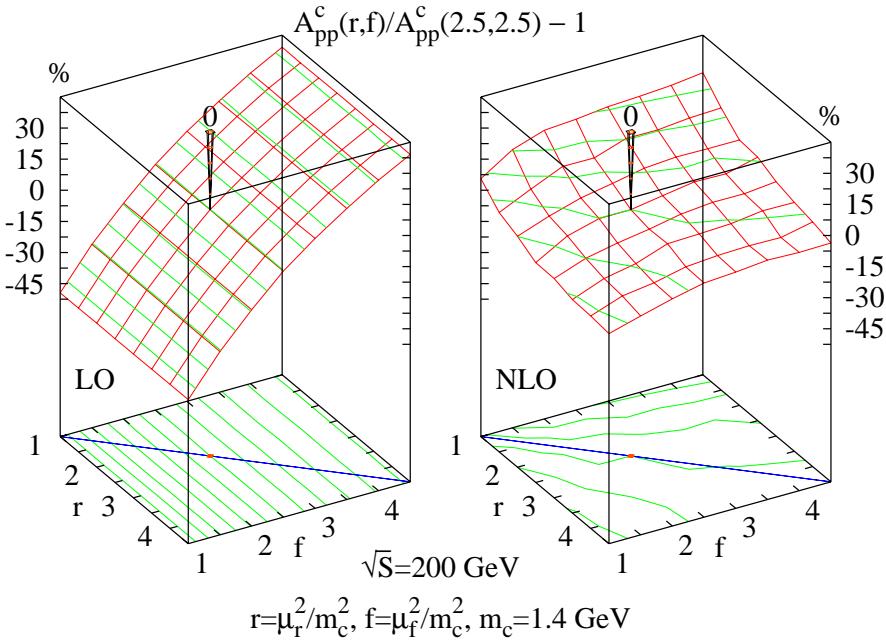
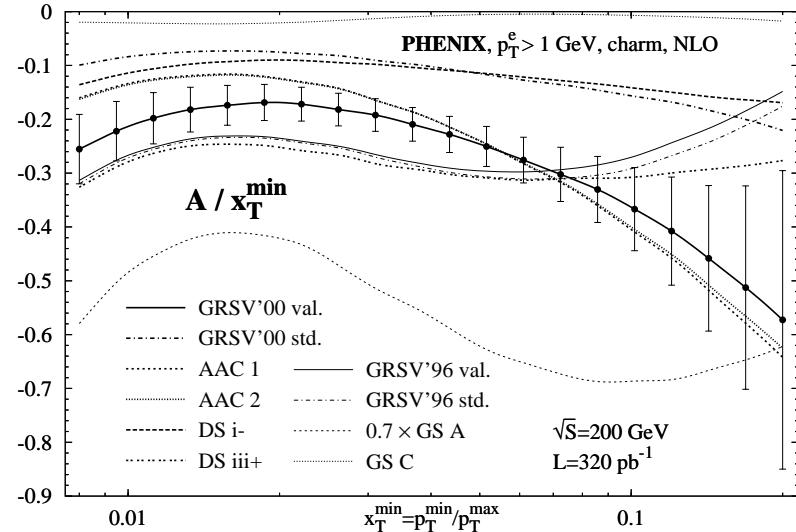
somewhat smaller  $x$  values accessible here

**complication:**  $c, b$  identified via their decay leptons

→ have to simulate HQ-decays in estimates of  $A_{LL}^Q$

$A_{LL}^c(p_T^{\min})$  for  :

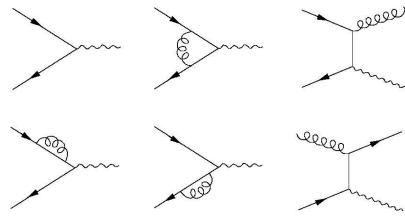
→ sensitivity to  $\Delta g$



$\mu_{r,f}$  dependence  
reduced in NLO

- $W^\pm$  boson production:

NLO: Weber; Kamal; Gehrmann; Smith et al.



**idea:** parity violating  $\rightarrow$  single-spin asymmetry sufficient

in LO one has e.g.

Leader, Sridhar; Bourrely, Soffer

$$A_L^{W^+}(y) = \frac{-\Delta u(x_a)\bar{d}(x_b) + \Delta \bar{d}(x_a)u(x_b)}{u(x_a)\bar{d}(x_b) + \bar{d}(x_a)u(x_b)} \text{ where } x_{a,b} = \frac{M_W}{\sqrt{S}}e^{\pm y}$$

$$\rightarrow y \text{ large} \& \begin{cases} \text{positive } (x_a > x_b) \\ \text{negative } (x_a < x_b) \end{cases} : A_L^{W^+} \simeq \begin{cases} -\Delta u(x_a)/u(x_a) \\ \Delta \bar{d}(x_a)/\bar{d}(x_a) \end{cases}$$

in practice:  $\Delta u$  is probed for  $x \gtrsim 0.2$  and  $\Delta \bar{d}$  for  $x \lesssim 0.12$

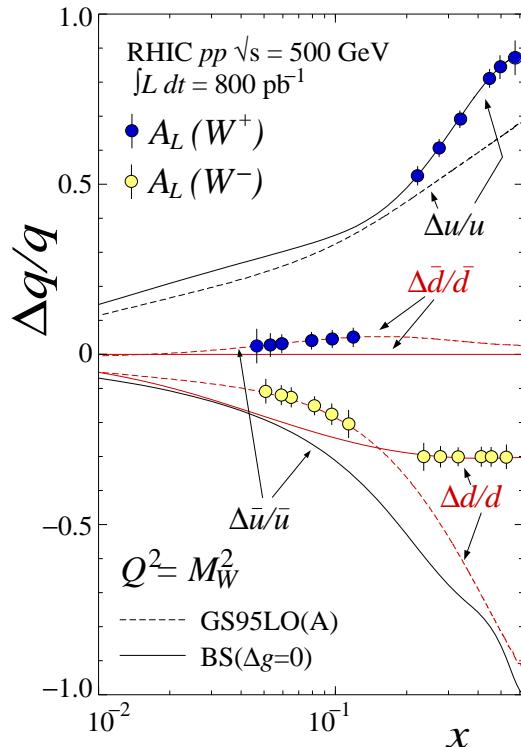
$A_L^{W^-}(y)$ :  $u \leftrightarrow d \rightarrow$  probes  $\Delta d/d$  and  $\Delta \bar{u}/\bar{u}$

**complication:** have to observe  $W^\pm$  via decay

but full NLO *lepton-level* MC is available

Nadolsky, Yuan

**expected sensitivity on  $\Delta q/q$  from  $A_L^{W^\pm}(y)$ :**

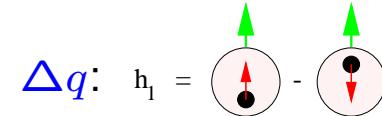


## Transverse Polarization

boosts and rotations do not commute → new object

**transversity density**  $\delta q(x, Q^2)$ : (sometimes also called  $h_1$ )

- “as partonic as”  $q$  and  $\Delta q$ :  $h_1 = \frac{1}{\sqrt{2}}(q + \Delta q)$



Ralston, Soper; Artru, Mekhfi; Jaffe, Ji

- reveals **chiral-odd** nature if described in helicity-basis

$$|\perp\rangle = \frac{1}{\sqrt{2}}(|+\rangle + |-\rangle) \text{ and } |\top\rangle = \frac{1}{\sqrt{2}}(|+\rangle - |-\rangle)$$

→ difficult to measure (fund. interactions preserve chirality)

→ no gluon density for spin 1/2 targets possible

Jaffe, Ji; Artru, Mekhfi; Ji

- simple non-singlet type evolution eqs.

NLO  $\delta P_{qq}$ : Vogelsang; Kumano, Miyama; Koike et al.

- vanishes at large scales (all  $\int_0^1 x^{n-1} \delta q(x, Q^2) dx$  decrease)
- constraint: Soffer; Sivers

$$|\delta q(x, \mu^2)| \leq \frac{1}{2} [q(x, \mu^2) + \Delta q(x, \mu^2)]$$

↑                      ↑                      ↑  
unknown              known (more or less)

- $\delta q$  could provide info on chiral-sym. breaking in QCD  
Collins; Jaffe

- tensor charge  $\sum_q \int_0^1 [\delta q - \delta \bar{q}](x, Q^2) dx \leftrightarrow$  neutron EDM  
Jaffe, Ji

## accessing $\delta q$ at RHIC:

chirality has to be flipped *twice* to access  $\delta q$

**double-spin    single-spin  
asymmetries**

$A_{\text{TT}}$                    $A_N$

**$A_{\text{TT}}$ :** ✓ clean, involves only  $\delta q$  as “unknowns” (but quadratic)

✗ small due to lack of gluons → exp. challenge

**$A_N$ :** ✗ requires *other unknown* chiral-odd fct. → involved

✓  $pp \uparrow \rightarrow \pi X$  sizable (E704, STAR)

### (I) double transverse spin asymmetries $A_{\text{TT}}$ :

on the menu: Drell-Yan, jets/hadrons, prompt-photons

general problem: ✗ “transversity” gluons

→ all asymmetries  $A_{\text{TT}} = d\delta\sigma/d\sigma$  strongly diluted by  
large number of gluon induced processes in  $d\sigma$

#### • Drell-Yan process

Ralston, Soper; Ji  
Cortes, Pire, Ralston; Artru, Mekhfi; Jaffe, Ji

✓ best candidate:  $q\bar{q} \rightarrow \mu^+\mu^-$  → no gluons in LO

✗ numerical study in NLO QCD revealed:

Martin, Schäfer, MS, Vogelsang

measurement suffers from limited muon acceptance

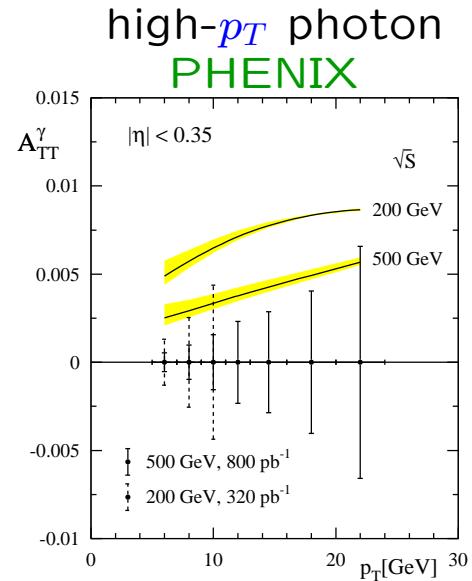
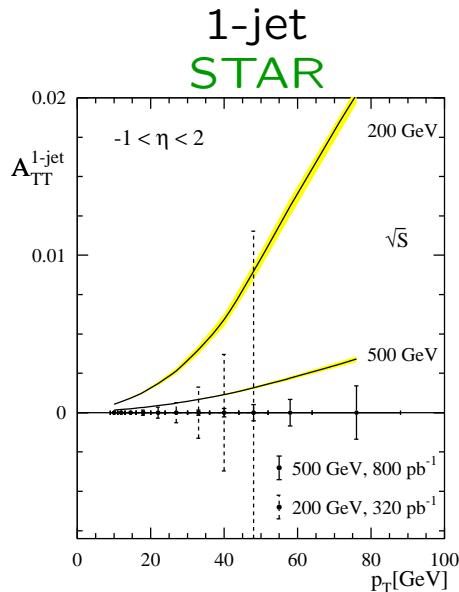
of :  $1.2 \leq |y_{\mu^\pm}| \leq 2.4$ ; need *both* muons!

→ conclusion on DY: tough, interesting with more lumi

- **high- $p_T$  jet and prompt photon production**

Hidaka et al.; Artru, Mekhfi; Ji; Jaffe, Saito; Soffer, MS, Vogelsang

- ✓ jets (and photons) copiously produced at colliders
- ✗ cannot avoid dilution of  $A_{TT}$  by gluon-induced processes

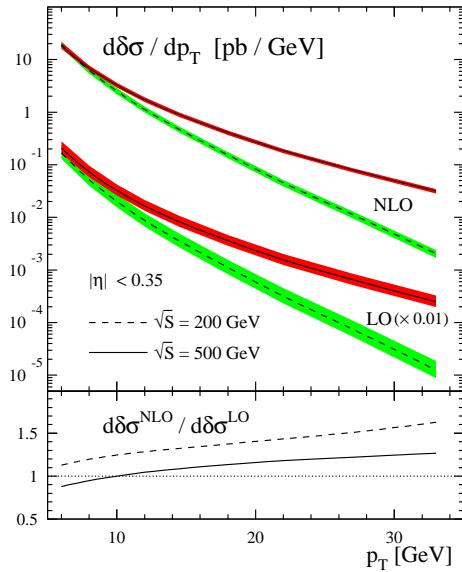


[upper bounds for  $A_{TT}$  in LO QCD; bands:  $\mu_F$  dependence]  
Soffer, MS, Vogelsang

**very recently:**

1<sup>st</sup> NLO computation of  
prompt photon production  
Mukherjee, MS, Vogelsang

[more involved due to  
extra spatial direction]



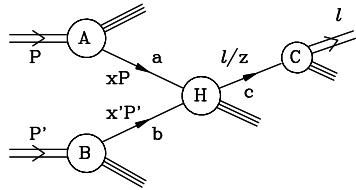
**general conclusions on  $A_{TT}$  at RHIC:**

- difficult, albeit not impossible task
- requires great exp. efforts [in particular, small syst. errors]
- large  $A_{TT}$ 's would constitute a new “spin surprise”

## (II) single transverse spin asymmetries $A_N$ :

exciting observable, goes back to the early days of spin

$A(p, \vec{s}_T) + B(p') \rightarrow C(l) + X$  with  $C$  '=' high- $p_T$   $\pi, \gamma, \dots$

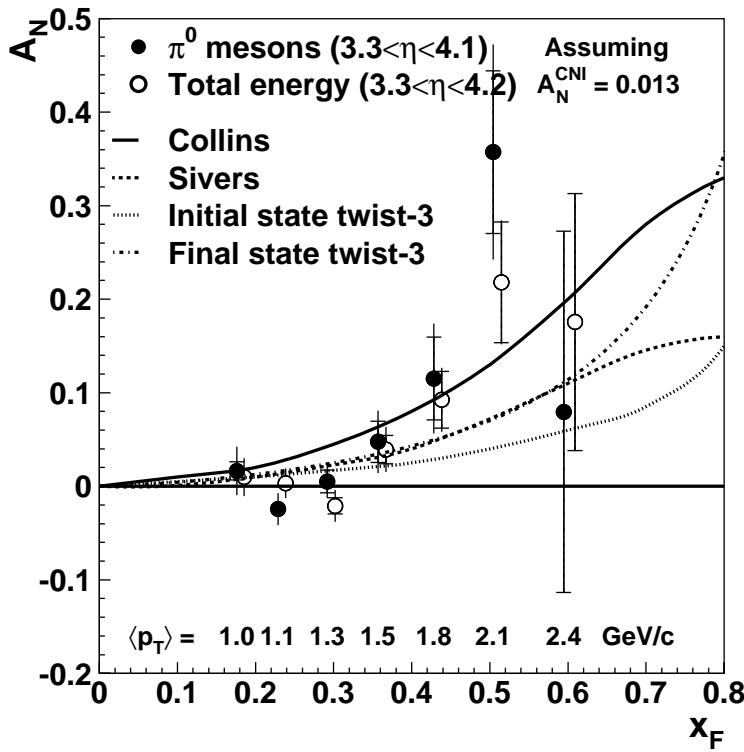
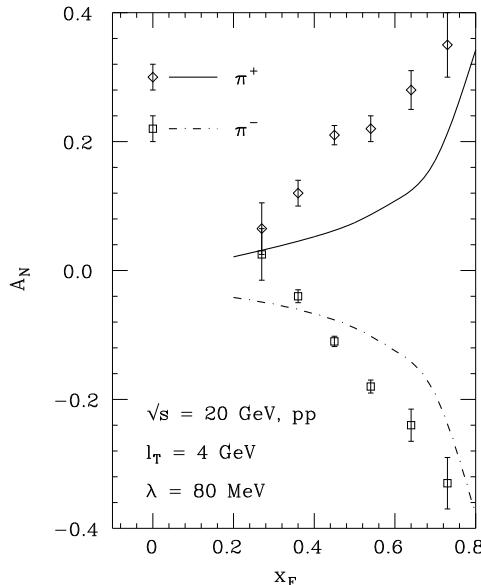


$$\text{measure: } A_N = \frac{\Delta_N \sigma}{\sigma} \equiv \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}}$$

“long-standing puzzle”:

large  $A_N$  for  $\pi$  and  $\Lambda$  production  
e.g. E704 ( $\sqrt{S} = 20 \text{ GeV}$ ) →

leading-twist pQCD:  $A_N = 0$



effect recently confirmed

also at  $\sqrt{S} = 200 \text{ GeV}$

by STAR

$A_N$  opens a window to a new class of power-suppr. “pdfs”

“twist-3” can arise from

**quark-gluon correlations    pdfs/ffs with intrinsic  $k_T$**

Qiu, Sterman; Koike et al.; Sivers; Collins; Boer; Leader et al.;

Teryaev et al. Anselmino et al.; Brodsky et al.;

Bacchetta et al.; ...

### problem:

very difficult to disentangle all these contributions

each mechanism on its own can account for E704 & STAR



progress requires combined efforts of  
several experiments and theory

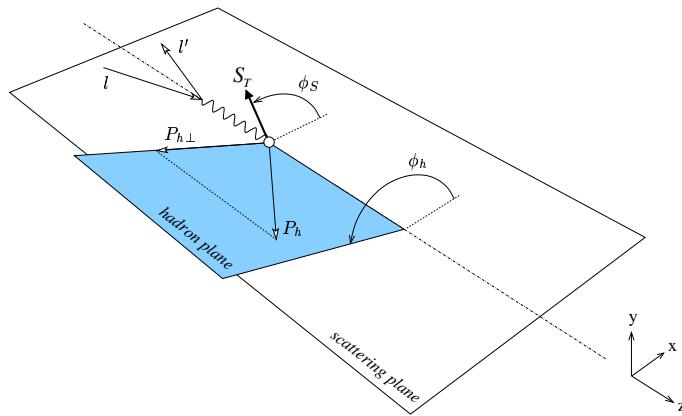
$pp$  (RHIC),  $e^+e^-$  (BELLE),  $lp$  (HERMES, COMPASS)

**intrinsic  $k_T$ :** a lot of theoretical activity recently

- mapping of all possible  $k_T$  distributions
- importance of gauge links  $\mathcal{P} \exp (ig \int_0^y dx^\mu A_\mu(x))$
- universality of distributions
- evolution, NLO, ...

→ Morning Session on Thursday

# very recent results on $ep^\uparrow \rightarrow e\pi X$ :



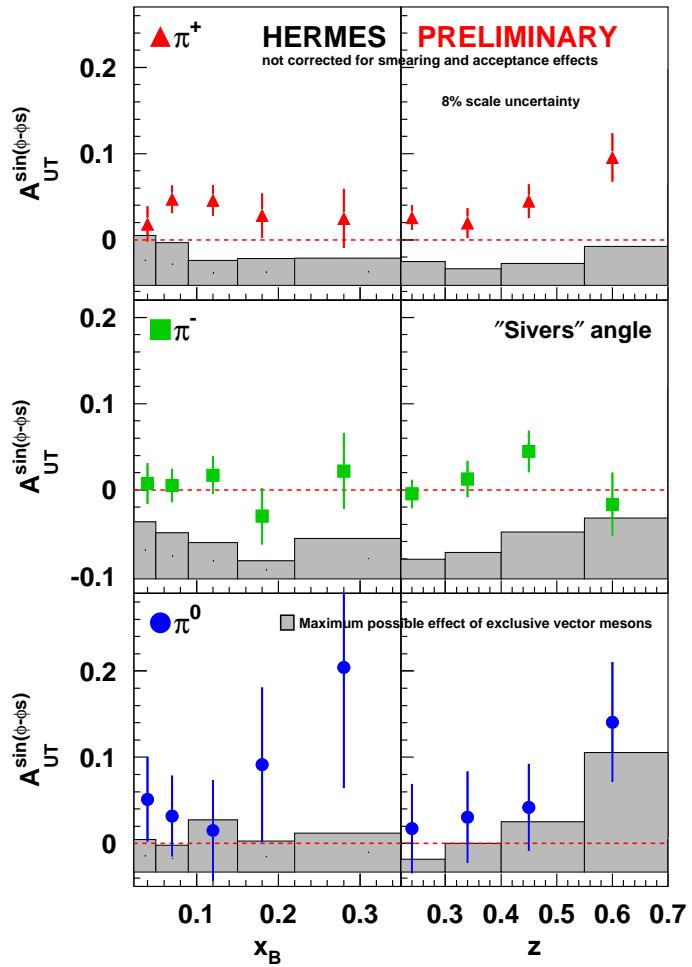
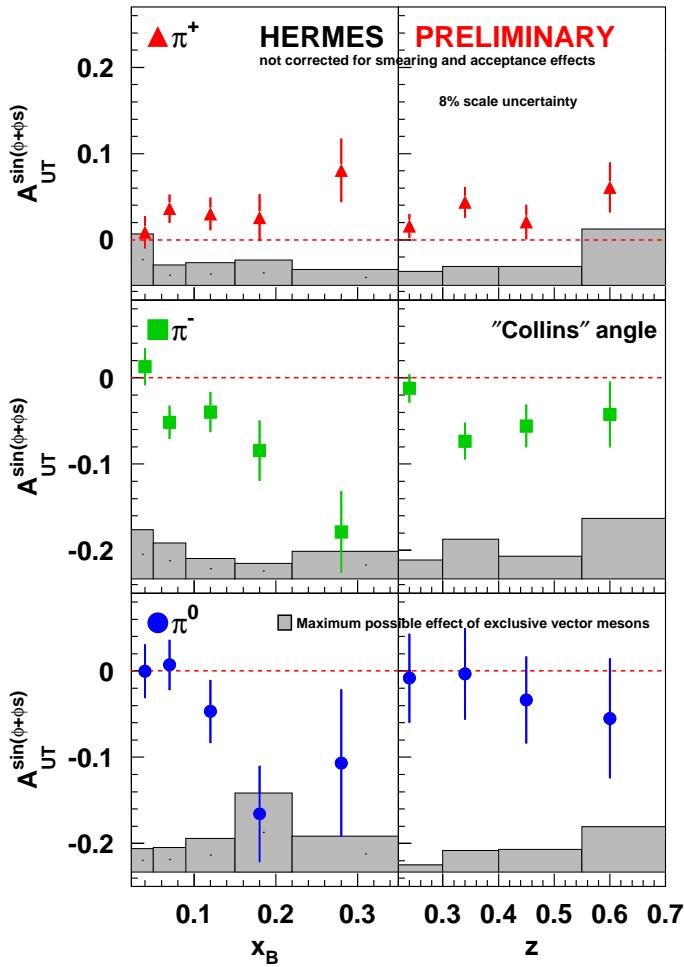
two leading contributions to  $A_{UT}$ :

$$\propto \sin(\phi_h + \phi_S) \delta q(x) H_1^{\perp,q}(z)$$

“Collins”

$$\propto \sin(\phi_h - \phi_S) f_{1T}^{\perp,q}(x) D_q(z)$$

“Sivers”



## Summary & Outlook

topics not touched:  $L_z$ , GPD's, lattice, ...

RHIC is the ideal playground to  
further test and learn much more  
about QCD spin interactions

first exciting results:  
QCD calculations will now be  
challenged by experiment

for many years to come RHIC will be  
the prime source of information on spin

next “spin surprise” is perhaps  
just round the corner

**very exciting times ahead of us**